

First-Principle Simulation of Particle Transport in the inversed-density-gradient profile

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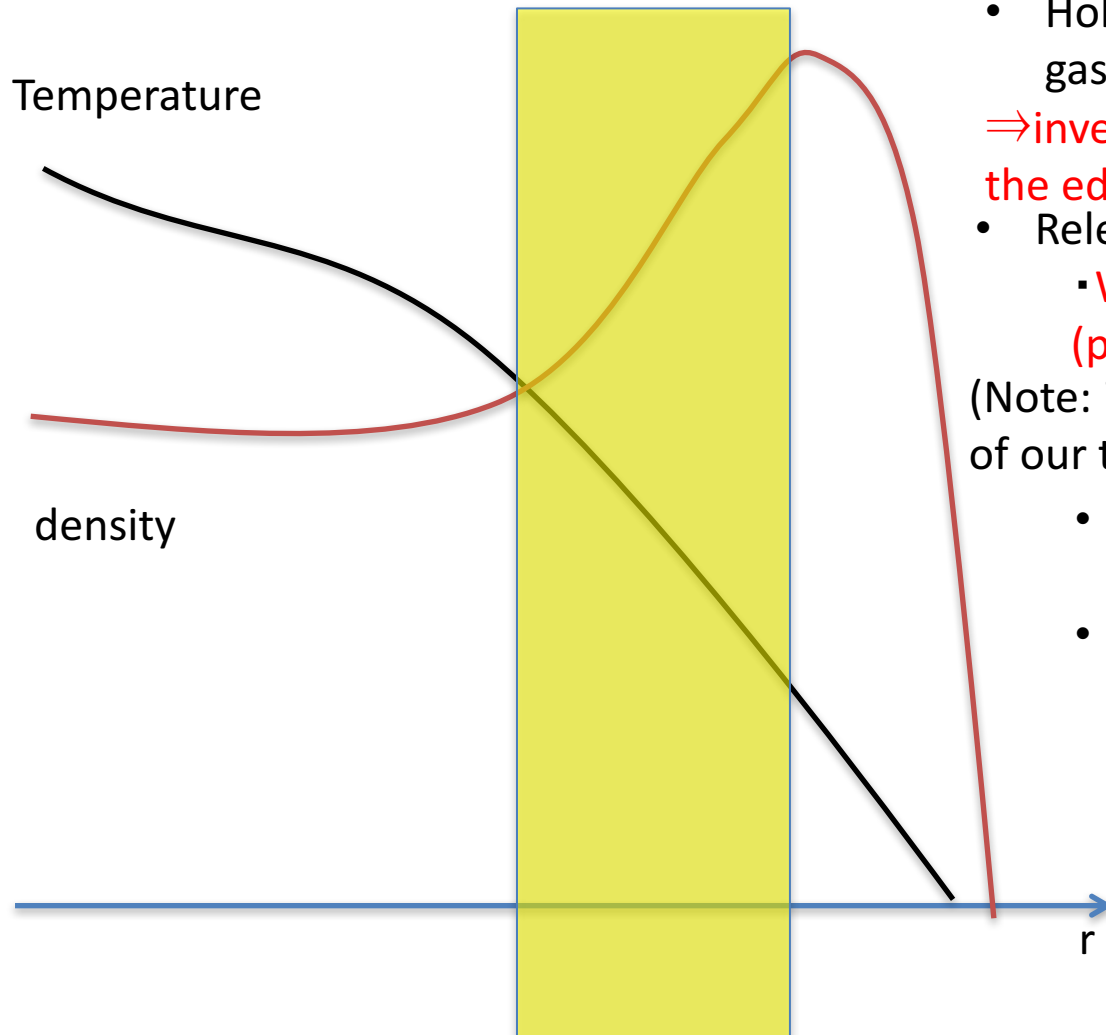
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- Summary

Inversed-density-gradient appears in fueling operations.



- Hollow density profile due pellet or gas-puff injection
 \Rightarrow **inversed density gradient appears in the edge region.**
- Relevant to
 - **What mechanism of the particle flux (pinch?)**

(Note: in this work, profile relaxation is out of our target, due to local approach.)

- Fueling physics contributes ITER/DEMO development.
- Works on validation of particle transport physics. [Angioni '09][Wan '10][Tangered '16]

We here undergo the first-principle simulation on the particle transport in the inversed-density gradient region.

dFEFI: delta- f gyrokinetic solver for ions and electrons. [B. Scott, PoP (2010)]

- delta- f Electromagnetic GK equation

$$\frac{\partial g}{\partial t} + \frac{c}{B_0} [(J_0 \psi_e), h]_{xy} - \frac{mz^2 + wB}{2e} \mathcal{K}(h) + \frac{B^s}{mB} [H_0, h]_{zs} = 0$$

$$g = \delta f + \frac{F^M}{T} e \frac{z}{c} J_0 A_{\parallel}, \quad h = \delta f + \frac{F^M}{T} e J_0 \phi, \quad H_0 = m \frac{z^2}{2} + wB,$$

$$\psi_e = J_0 \left(\phi - \frac{z}{c} A_{\parallel} \right).$$

$$F^M = n(2\pi T/m)^{-3/2} \exp(-H_0/T),$$

Polarization equation:

$$\sum_{\text{sp}} \int d\mathcal{W} \left[e J_0 g + e^2 \frac{F^M}{T} (J_0^2 - 1) \phi \right] = 0$$

Induction equation:

$$\nabla_{\perp}^2 A_{\parallel} + \sum_{\text{sp}} \frac{4\pi}{c} \int d\mathcal{W} \left[e z J_0 g - \frac{e^2}{c} z^2 \frac{F^M}{T} J_0^2 A_{\parallel} \right] = 0$$

- Field-aligned coordinate
 - Shifted metric
 - Fixed boundary on radius
 - Local code, but remains globality on the boundary conditions.
- (x: radial, y: binormal, s: magnetic field line)

Benchmark on cyclone-base case(CBC) parameter

core plasma:

$a/R = 0.184$, $L_{\perp}/R = 0.145$, $\min(k_y \rho_s) = 0.025$. $T_i/T_e = 1$, $qR/L_{\perp} = 9.67$,

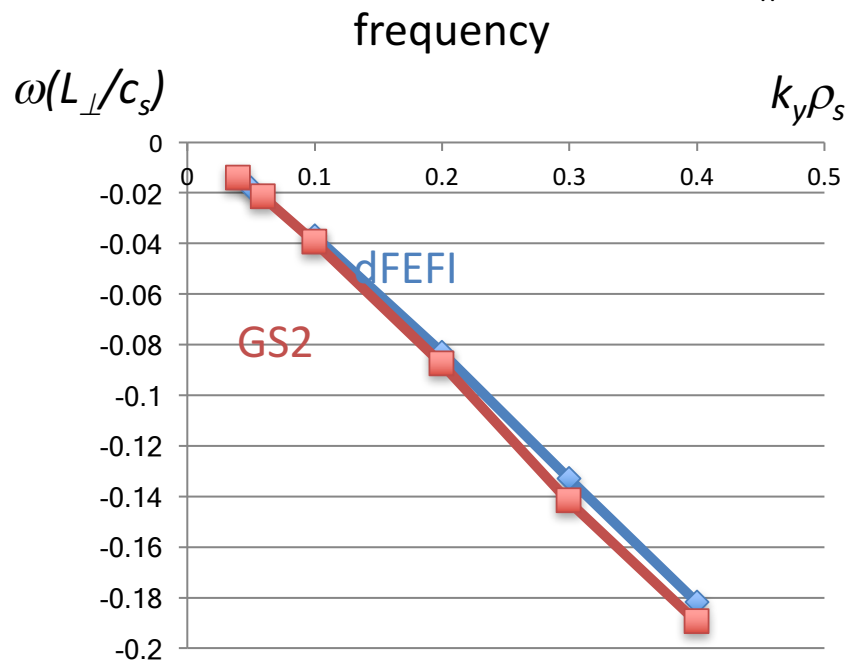
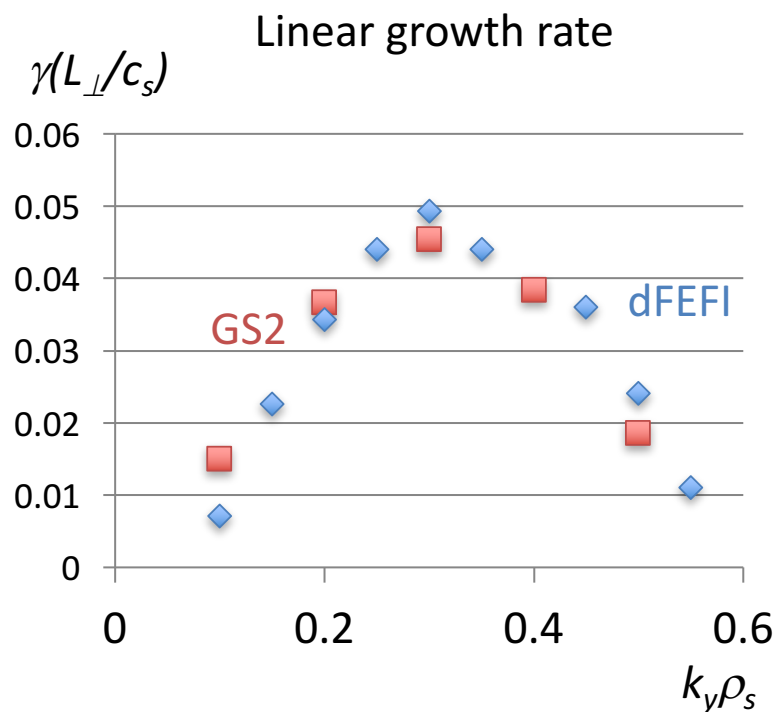
$\beta_e (qR/L_{\perp})^2 = 1e-3$, $rq'/q = 1.14$, $L_{\perp}/L_n = 0.321$, $L_{\perp}/L_{te} = L_{\perp}/L_{Ti} = 1.0$

$L_y/L_x = 4.0$, $(n_x, n_y, n_s, n_z, n_w) = (128, 128, 32, 48, 16)$

$R/L_{Ti} = 6.9$

$R/L_{Te} = 6.9$

$R/L_n = 2.2$

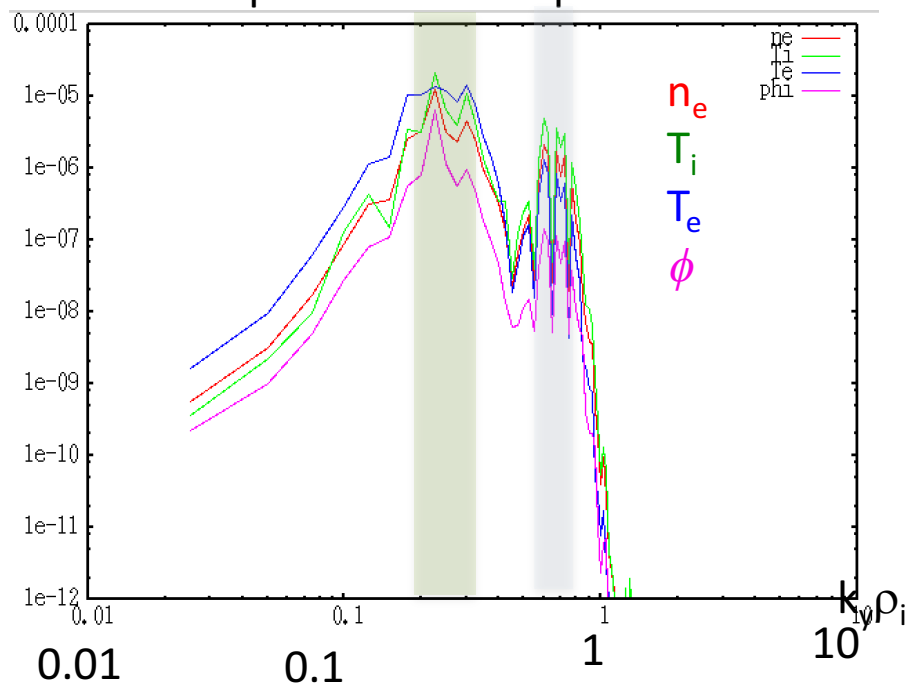


Linear calculation exhibits two modes with ion and electron direction rotation

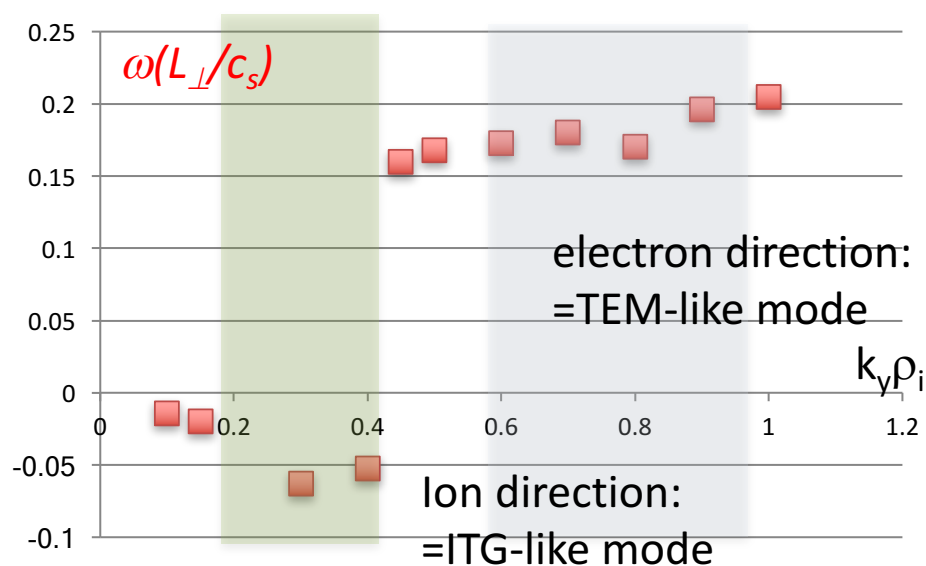
Edge plasma parameters:

$R_0=165\text{cm}$, $n_e=2.0\times 10^{13}\text{ cm}^{-3}$, $T_e=100\text{ eV}$, $L_n=-7\text{cm}$, $L_{Ti}=L_{Te}=3.5\text{cm}$, $B=2.5\text{ T}$, $T_i/T_e=1$, $a/R=0.303$,
 $L_\perp/R=0.0212$, $q=3.5$, Normalized beta= 6.44×10^{-5} , $v_i(L_\perp/c_s)=0.00956$, $v_e(L_\perp/c_s)=0.823$
 $(n_x,n_y,n_z,n_w)=(32,128,32,32,16)$, $L_x/L_y=1.0$, $\Delta t=0.005$

Power spectra in linear phase

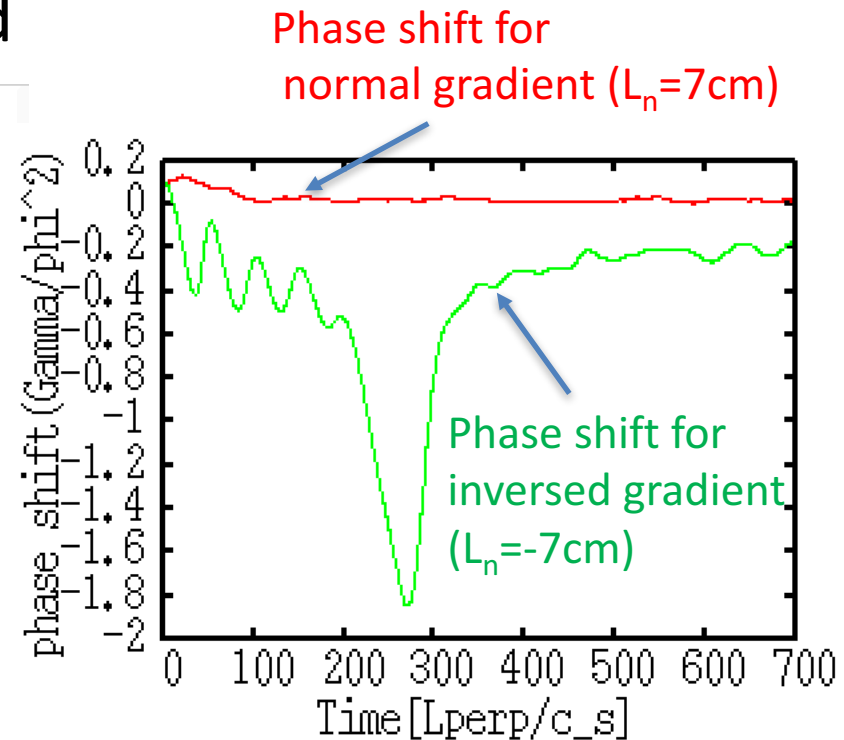
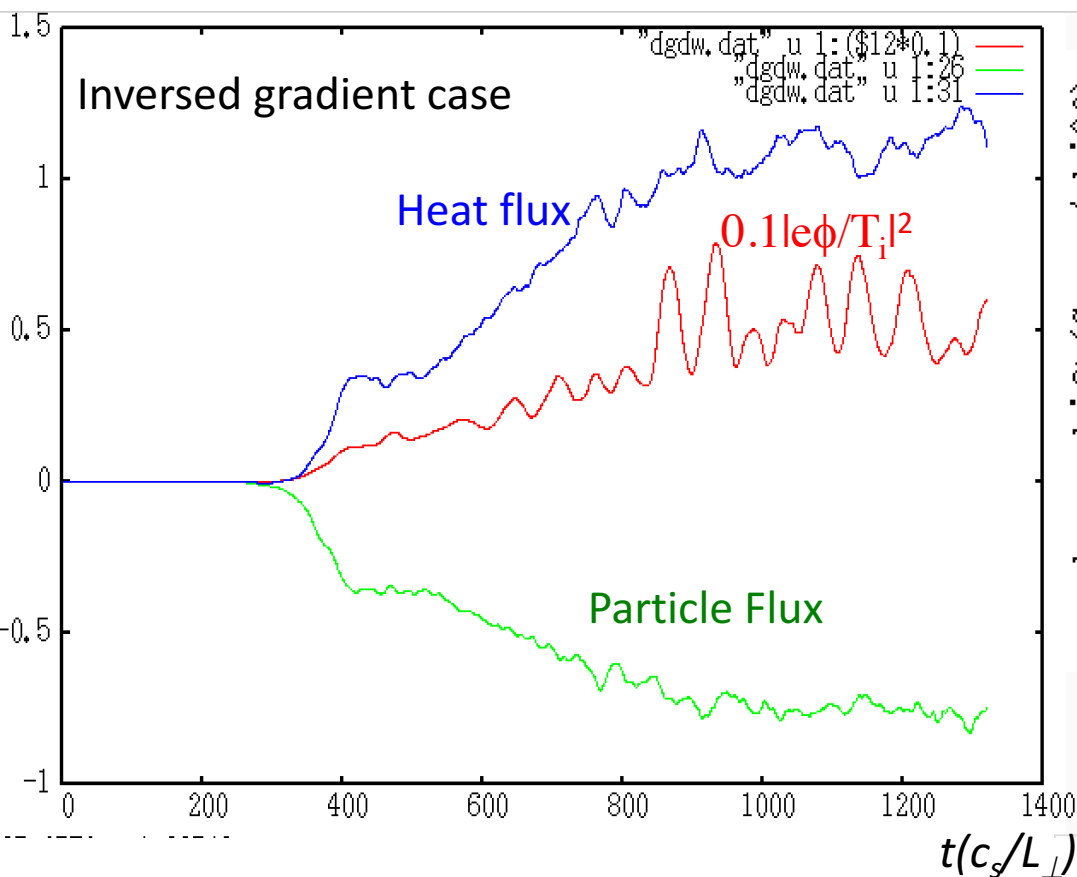


We observe ITG-like and TEM-like modes growing.



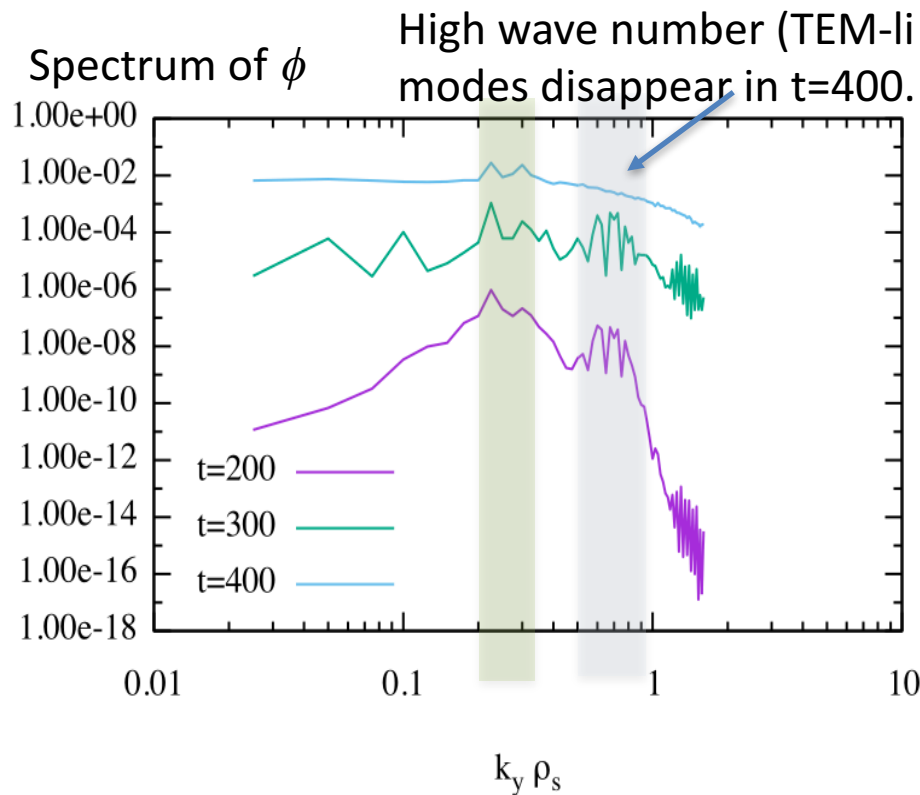
Time evolution of Particle Flux

⇒ Inward, with phase-shifted

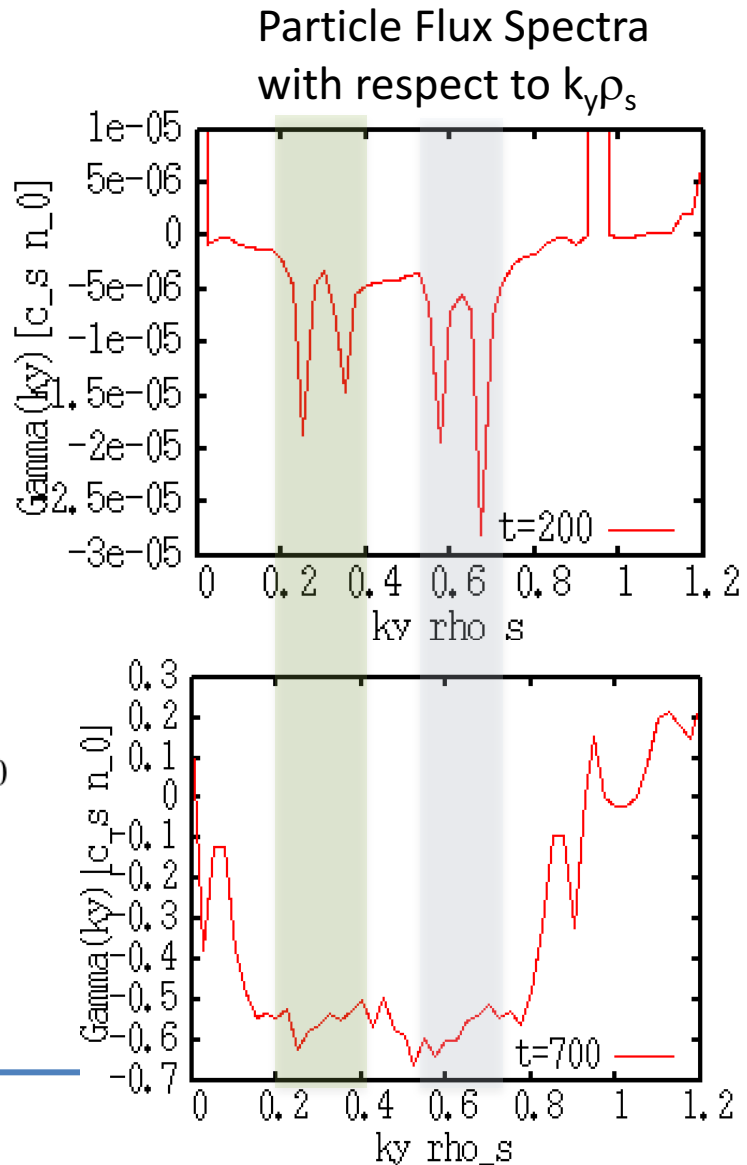


For the inversed-density-gradient case, we expect larger phase shifts in the density-potentials.

Spectral analyses on electrostatic potential and particle flux.



- In the nonlinear phase, the lower wave number mode survives.
- Both peaks contribute to inward particle fluxes. In the nonlinear phase, modes with $0.1 < k_y \rho_s < 0.8$ contribute to the inward flux.



Candidate? Ion-Mixing-Mode [Coppi '78 PRL]

- Slab fluid Model

- Mode dispersion

$$1 + iA + k_{\parallel}^2 \frac{T_e}{m_i} \frac{\omega_{Ti}}{\omega^3} - \frac{\omega_{*e}}{\omega} = 0$$

- Non-trivial mode:

$$\omega = -[(k_{\parallel}^2 T_e / m_i) \omega_{Ti} / (1 + A^2)]^{1/3} (1 - iA)^{1/3}$$

ITG-like mode

Electron adiabaticity is affected by the electron thermal force.

$$\frac{\tilde{n}_e}{n} = \frac{e\tilde{\phi}}{T_e} \left[1 + (1 + \alpha_T) \frac{i}{\omega_{\chi}} (\omega - \omega_{*e} + \frac{3}{2} \eta_e \omega_{*e}) \right]$$

$$\omega_{\chi} \equiv \hat{\chi}_e k_{\parallel}^2 T_e / (m_e v_e)$$

Assumption: $\omega \sim \omega_{*e} < \omega_{\chi}$ \longrightarrow Diagonal effects (prop. to grad n) are cancelled, remaining off-diagonal (prop. to grad Te) term.

Particle Flux: $\Gamma = \langle \tilde{n} \tilde{v}_{E_x} \rangle = -(2c / B) \text{Im}(\sum_k k_y \tilde{\phi}_k \tilde{n}_k)$

$$= 3(1 + \alpha_T) D_B^i n \sum_k k_y^2 \left| \frac{e\tilde{\phi}}{T_e} \right|^2 \frac{1}{\omega_{\chi}} \frac{c T_e}{e B} \frac{d \ln T_e}{dx}$$

However, Is this assumption correct ?

Estimate from simulation results ~ -0.73

Validation of Ion-Mixing-Mode

$$\omega \sim \omega_{*e} < \omega_\chi \text{ Is correct?}$$

$$\frac{\tilde{n}_e}{n} = \frac{e\tilde{\phi}}{T_e} \left[1 + (1 + \alpha_T) \frac{i}{\omega_\chi} (\omega - \omega_{*e} + \frac{3}{2} \eta_e \omega_{*e}) \right]$$

$$\omega(k_y = 0.1) \sim -0.063(c_s / L_\perp)$$

$$\omega_{*e} = \frac{k_y c T_e}{eB} \frac{d \ln n}{dr} \sim -0.1 (c_s / L_\perp)$$

$$\omega_\chi = \hat{\chi}_e k_\parallel^2 T_e / m_e v_e \sim 0.7(c_s / L_\perp) > \omega_{*e}$$

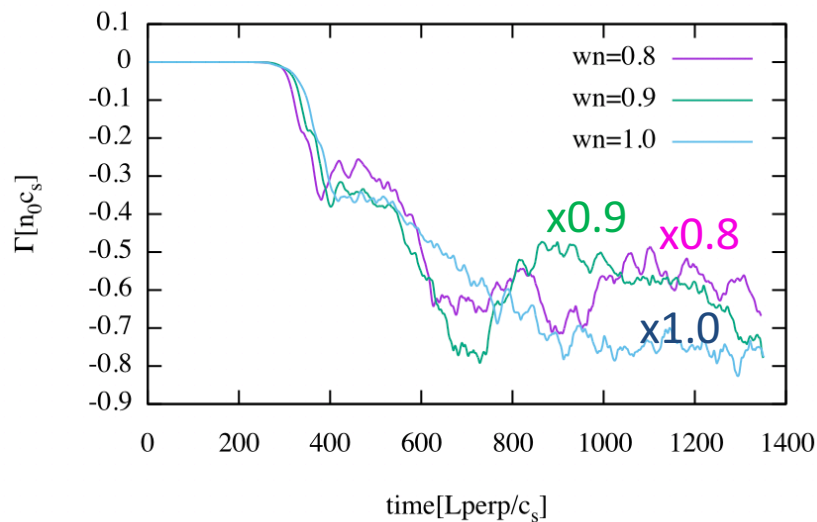
$$A = \frac{3}{2} \eta_e \frac{\omega_{*e}}{\omega_\chi} (1 + \alpha_T) \sim 0.73 > \sqrt{\frac{m_e}{m_i}} \rightarrow \text{This mode is enough large compared with collisional effects.}$$

dTe/dr dependency

Measures degree of the cancellation.

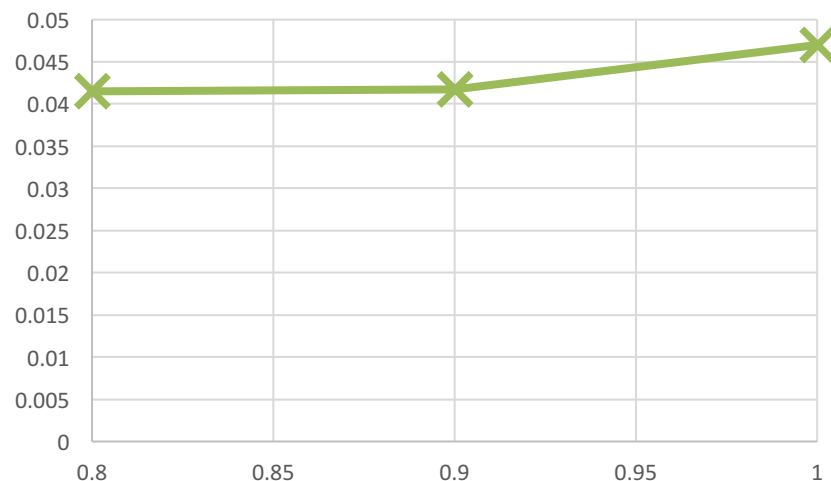
Parameter Scan(1) n scan

$\Gamma(n_0 c_s)$



Growth rate

γ

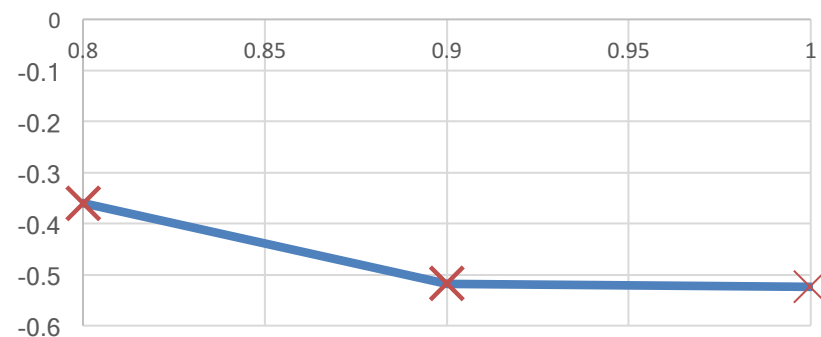


- ∇n scan on saturated particle flux is almost similar within fluctuations.
 - Diffusive part is not significant
- Higher ∇n gives (slightly) higher growth rate.

In the IMM theory, a diffusive part on particle flux is not significant, since $\omega \sim \omega_{*e}$ is satisfied.

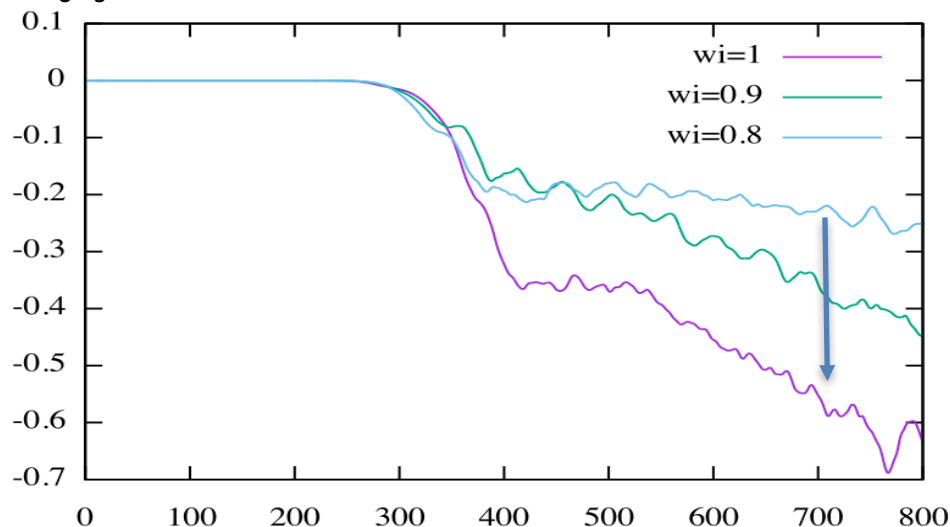
$\delta(=\langle nv \rangle / \phi^2)$

$(\nabla n) / (\nabla n)_{\text{ref}}$



Parameter scan(2) Grad Ti Scan

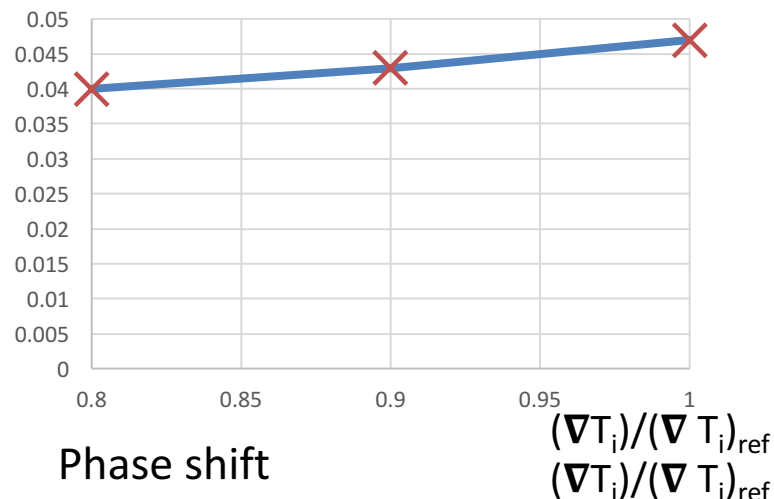
$\Gamma(n_0 c_s)$ Grad Ti scan (x1.0, x0.9, x0.8)



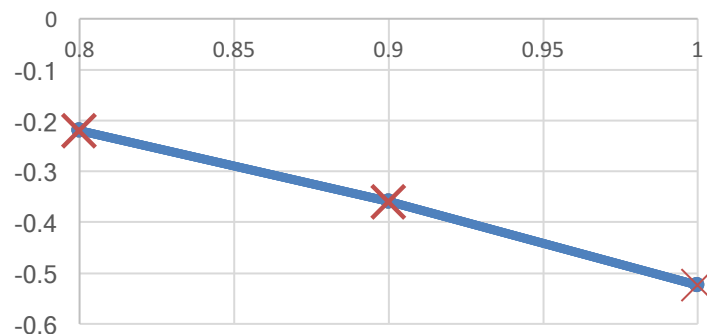
- Higher ∇T_i gives higher saturated inward particle flux.
- Higher ∇T_i gives higher growth rates, related to ITG mode.
- Higher ∇T_i gives drastically higher phase shift ($\delta = \Gamma / \langle \phi_2 \rangle$).

Higher $\nabla T_i \Rightarrow$ higher Inward flux, originated from phase shift, as well as increase of growth rates

Growth rate γ



$\delta (= \langle nv \rangle / \phi^2)$



summary

- We have simulated the turbulence in case that density gradient inversed. In linear phase, two modes appear; one is ITG-like, and the other is TEM-like modes.
- The TEM-like modes disappears in the nonlinear phase.
- The ITG-like mode may be identified as the ion-mixing-mode. Note that the IMM is unique to the inversed-density-gradient case.
- In future work, toroidal IMM model should be derived for further understanding of the simulation results.

Implications:

- Further physics should enter in hollow density profile.
- Electron temperature is a key for the case, in that electron thermal fluctuation can modulate the phase shift.